

# Dissipation and information in stochastic processes

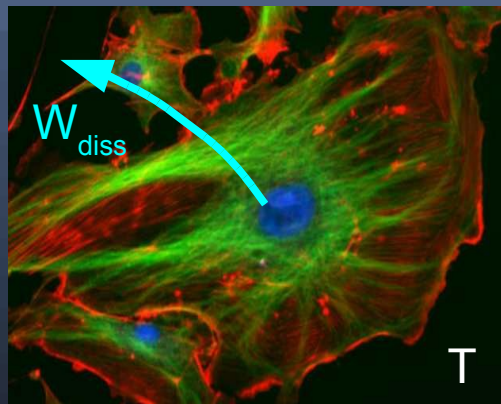
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Universidad Complutense de Madrid

GISC Workshop '10. February 19<sup>th</sup> 2010. Madrid (Spain).

# Dissipation and Irreversibility

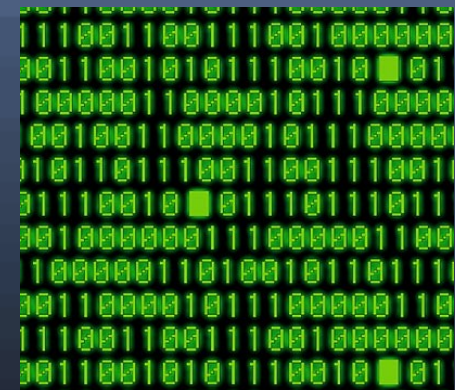
Dissipation-irreversibility relationship in non-equilibrium stochastic processes :



Energy dissipation



Fluctuation Theorems



Information  
Irreversibility revealed in the data

# Fluctuation Theorems (FT)

J.M.R. Parrondo, C.V. Den Broeck and R. Kawai, New Journal of Physics **11**, 073008 (2009).

$$\langle \dot{W}_{diss} \rangle = \langle \dot{W} \rangle - \Delta \dot{F} = T \langle \dot{S} \rangle = \lim_{t \rightarrow \infty} \frac{kT}{t} D \left[ p \left( \{x(\tau)\}_{\tau=0}^t \right) \middle| \middle| p \left( \{x(t-\tau)\}_{\tau=0}^t \right) \right]$$

Average dissipation of the process  
Entropy production

Relative entropy between forward and backward processes

D = Relative Entropy (Kullback-Leibler distance)

# Relative Entropy or Kullback–Leibler distance

$$D(p||q) = \sum_i p_i \log \left( \frac{p_i}{q_i} \right)$$

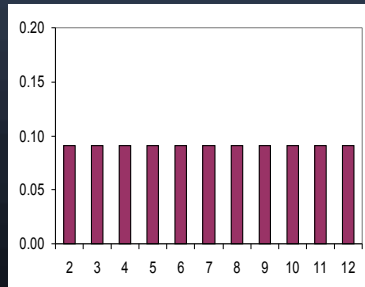
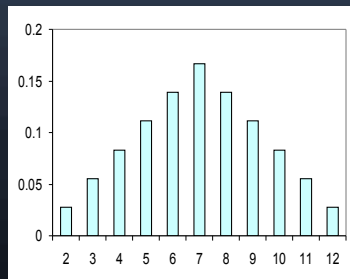
5,7,12,9,6,2,5,2,9,2,12,8,10,4,....



OR



?



Stein's Lemma:

The probability of incorrectly guessing  $p$  from  $n$  data actually distributed as  $q$  is (asympt.):

$$2^{-nD(p||q)}$$

$$D(\text{Dice} || \text{Ball Machine}) = 0.185$$

$$D(\text{Ball Machine} || \text{Dice}) = 0.220$$

# Fluctuation Theorems (FT)

$$\langle \dot{W}_{diss} \rangle = \langle \dot{W} \rangle - \Delta \dot{F} = T \langle \dot{S} \rangle = \lim_{t \rightarrow \infty} \frac{kT}{t} D \left[ p \left( \{x(\tau)\}_{\tau=0}^t \right) \middle| \middle| p \left( \{x(t-\tau)\}_{\tau=0}^t \right) \right]$$

Stochastic discrete processes in stationary regime:  $(x_1, x_2, \dots, x_n)$

$$D_k(p_F || p_B) = \sum_{x_1, \dots, x_k} p(x_1, \dots, x_k) \log \frac{p(x_1, \dots, x_k)}{p(x_k, \dots, x_1)}$$

$$\frac{\langle \dot{S} \rangle}{k} = d(p_F || p_B) = \lim_{n \rightarrow \infty} \frac{1}{n} D_n(p_F || p_B)$$

partial information

$$\frac{\langle \dot{S} \rangle}{k} \geq d(p_F || p_B)$$

Physics  
(Average)

Time series  
(single trajectory)

Even ignoring some physical details of the system we can estimate its dissipation !

# Relative entropy estimators: Brute force

$$\frac{\langle \dot{S} \rangle}{k} \geq d(p_F || p_B)$$

Problem: Calculating D from a single trajectory  
2,3,2,1,0,2,4,5,3,4,5,0,2,1,0,2,1,2,1, ...

1<sup>st</sup> estimator: Brute-force counting of k-th length strings

$$D_n(p_F || p_B) \equiv \sum_{x_1, \dots, x_n} p(x_1, \dots, x_n) \log \frac{p(x_1, \dots, x_n)}{p(x_n, \dots, x_1)}$$

$$d_k \equiv D_k - D_{k-1}$$

$$d(p_F || p_B) = \lim_{n \rightarrow \infty} \frac{1}{n} D_n(p_F || p_B) = \lim_{k \rightarrow \infty} d_k$$

k-th order Markov chains:

$$d = d_{k+1}$$

# Relative entropy estimators: Ziv-Merhav

2<sup>nd</sup> estimator: Based on Lempel-Ziv compression

$$D(p \parallel q) \sim \text{🔧} (p \parallel q)$$

$$d_{ZM}(\mathbf{x}) \equiv \frac{1}{n} [c(\mathbf{x}|\mathbf{x}') \log n - c(\mathbf{x}) \log c(\mathbf{x})]$$

$c(\mathbf{x})$  = Compression length of  $\mathbf{x}$  using the Lempel-Ziv incremental parsing

$c(\mathbf{x}|\mathbf{x}')$  = Cross parsing length of  $\mathbf{x}$  with respect to its reverse  $\mathbf{x}'$

Tends to relative entropy asymptotically.

# Relative entropy estimators: Ziv-Merhav

$c(z)$  , compression length of  $z$

$$z = ( 0 \ 1 \ 1 \ 1 \ 1 \ 0 \ 0 \ 0 \ 1 \ 1 \ 0 )$$

$$z = ( 0 \mid 1 \mid 11 \mid 10 \mid 00 \mid 110 ) \quad c(z) = 6$$

$c(z|x)$  , Cross parsing length of  $z$  with respect to  $x$

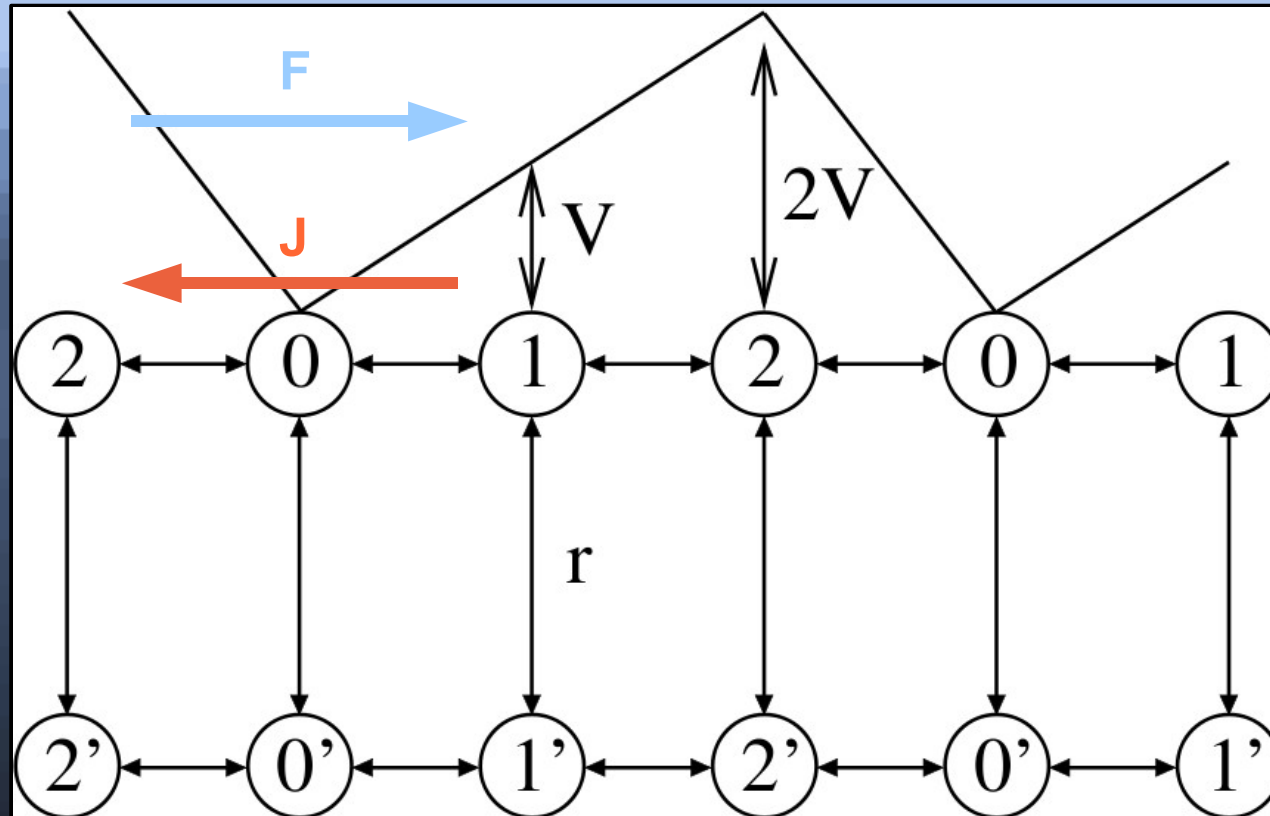
$$z = ( 0 \ 1 \ 1 \ 1 \ 1 \ 0 \ 0 \ 0 \ 1 \ 1 \ 0 )$$

$$x = ( 1 \ 0 \ 0 \ 1 \ 0 \ 1 \ 0 \ 0 \ 1 \ 1 \ 0 )$$

$$z = ( 011 \mid 110 \mid 00110 ) \quad c(z|x) = 3$$



# Discrete ratchet

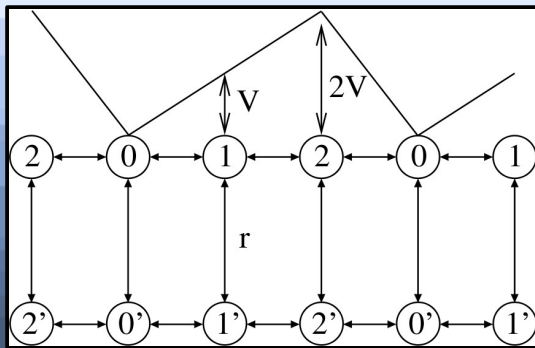


All rates obey detailed balance **except** switches → Irreversibility → **Dissipation**

$$\frac{k_{0 \rightarrow 1}}{k_{1 \rightarrow 0}} = e^{-\beta V} \cdot e^{\beta FL}$$

$$\frac{k_{1 \rightarrow 1'}}{k_{1' \rightarrow 1}} = \frac{r}{r} = 1 \neq e^{\beta V}$$

# Discrete ratchet: full information



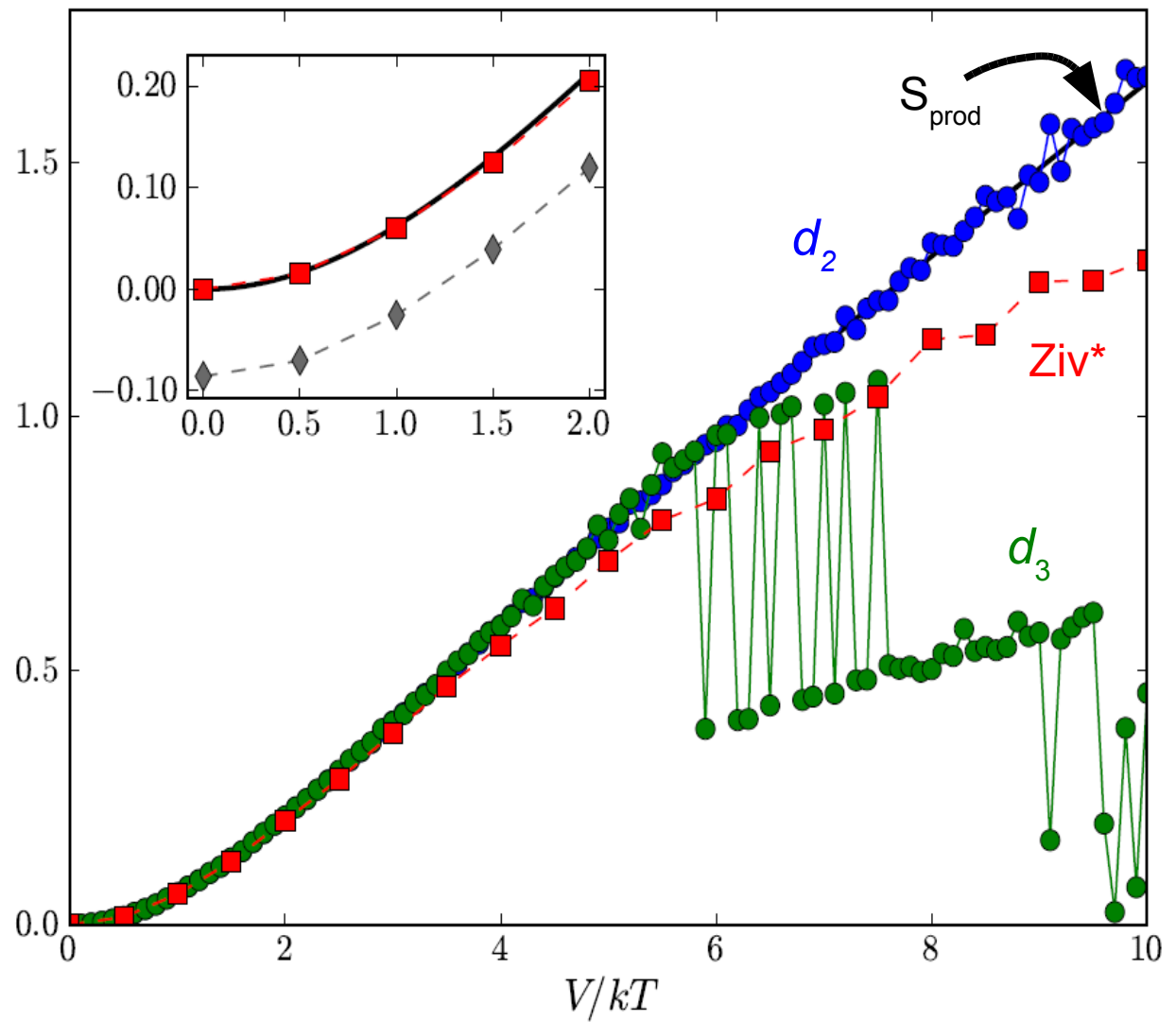
$F=0$   
 $r$  finite

Estimate dissipation  
with single trajectories

**Full Information**

$x=\{0,1,2\}$

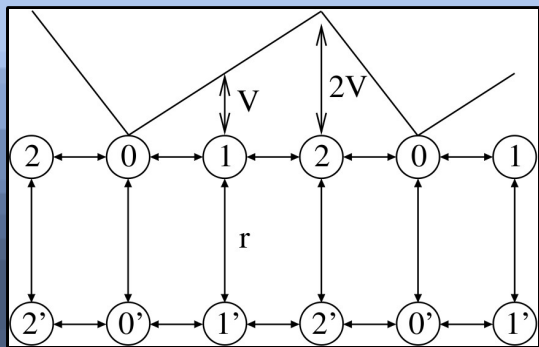
$y=\{0,1\}$



$$\frac{\langle \dot{S} \rangle}{k} \equiv d_2 = d_3 = d_4 = \dots$$

Markovian

# Discrete ratchet: partial information

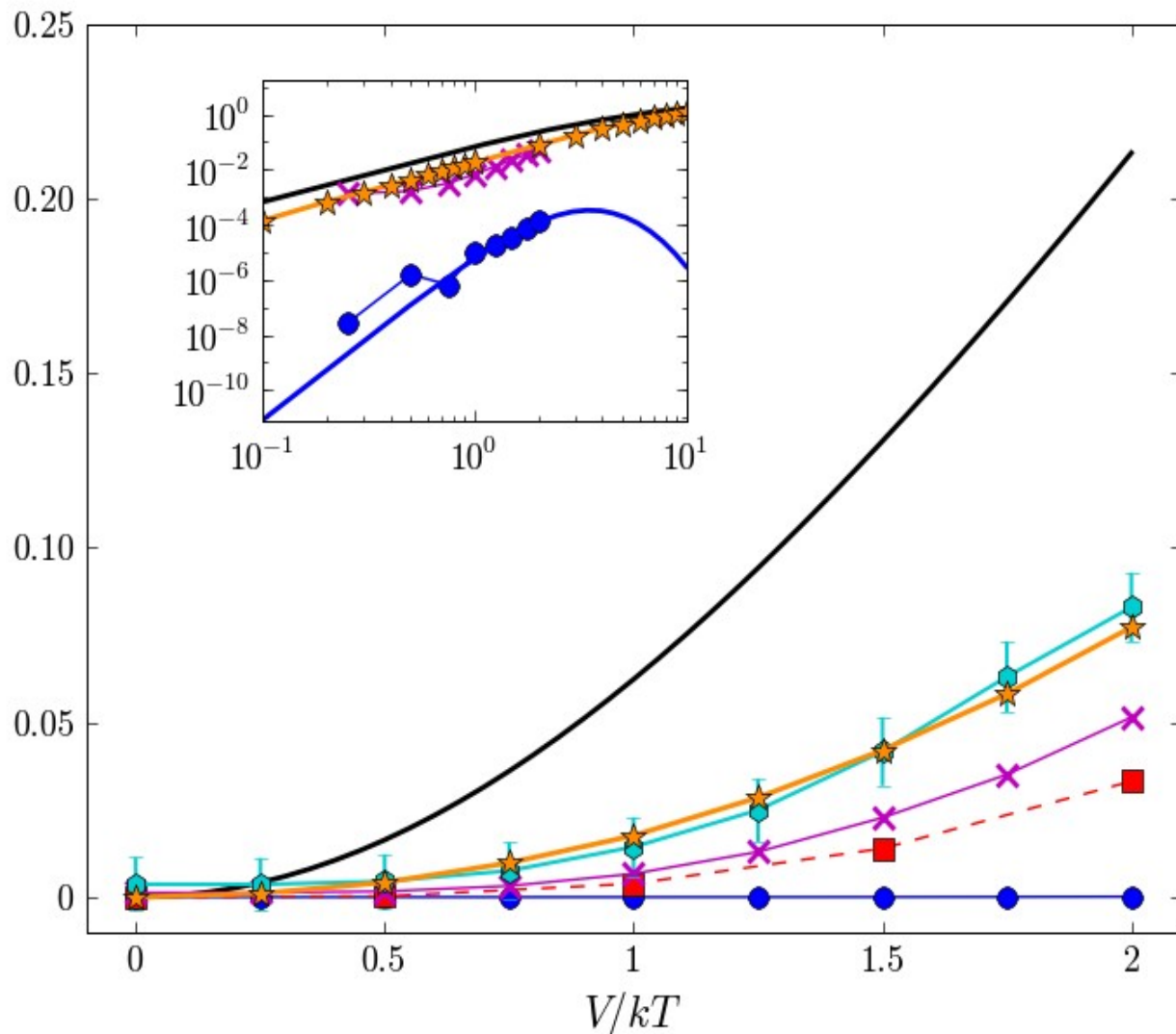


$F=0$   
 $r$  finite

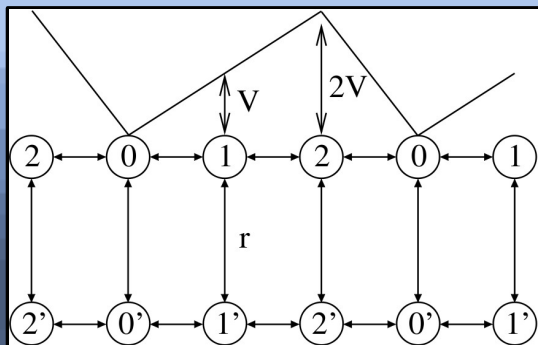
Partial Information  
 $x=\{0,1,2\}$

$$d_k \equiv D_k - D_{k-1}$$

$$D_k = \sum_{x_1, \dots, x_k} p(x_1, \dots, x_k) \log \frac{p(x_1, \dots, x_k)}{p(x_k, \dots, x_1)}$$



# Discrete ratchet: partial information



$F=0$   
 $r$  infinite

Partial Information  
 $x=\{0,1,2\}$

$$J_{ij} \equiv p_{ij} - p_{ji}$$

currents

$$J_{ij} \ll p_{ij}$$

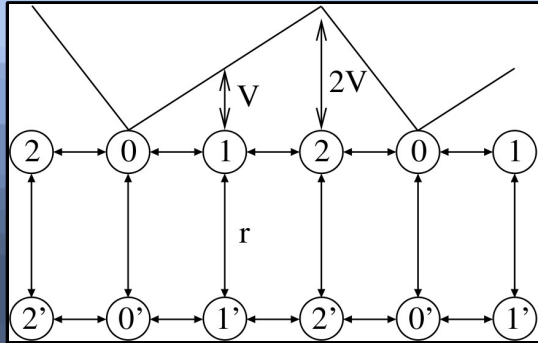
small current limit

$$d_2 \simeq \sum_{ij} \frac{(J_{ij})^2}{2p_{ij}} + O(J^3) \simeq \sum_{i<j} \frac{(J_{ij})^2}{p_{ij}} + O(J^3)$$

$$d_2 = d_3 = d_4 = \dots = d = 0 < \frac{\langle \dot{S} \rangle}{k}$$

Markovianity is not enough!

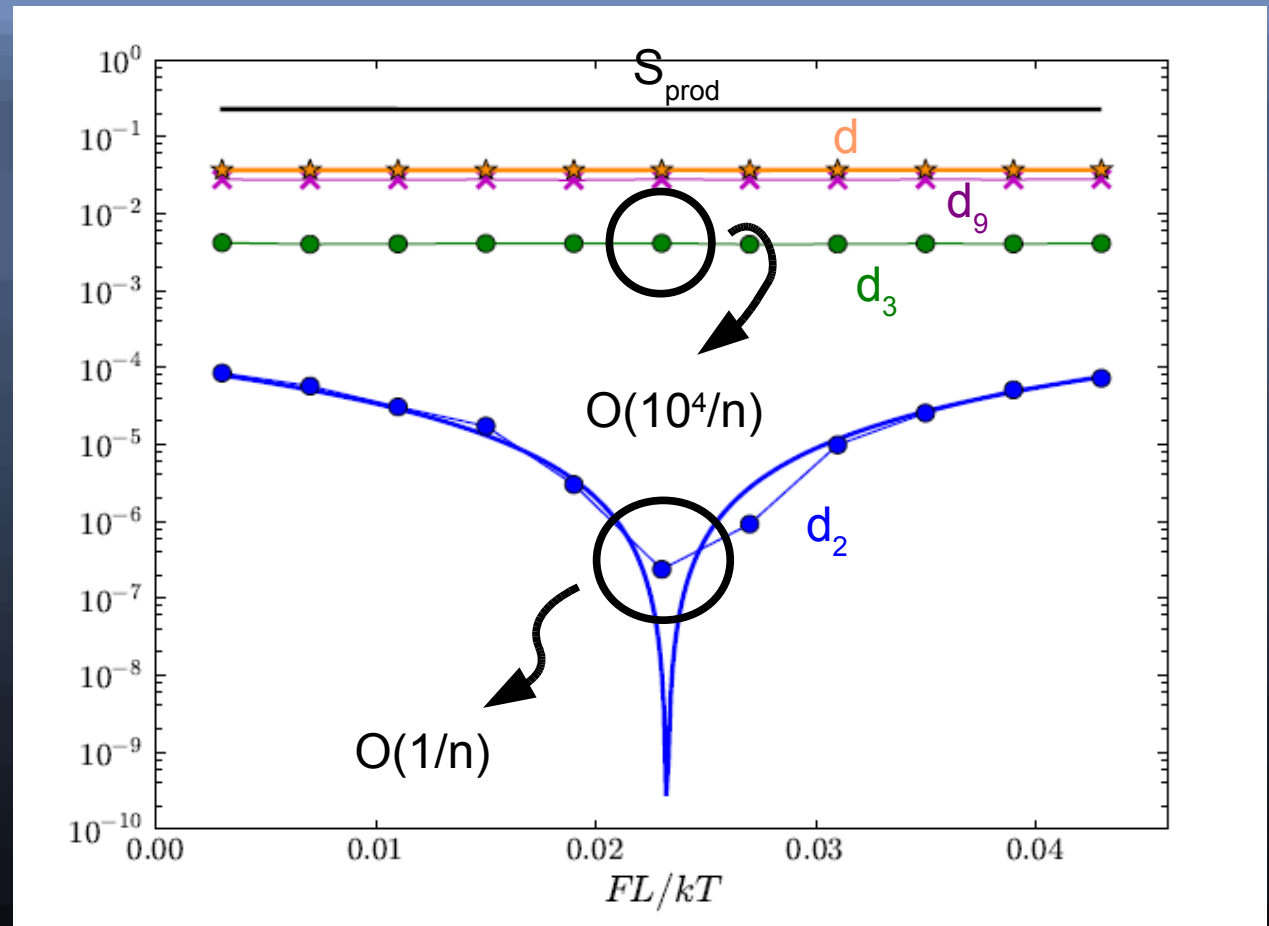
# Discrete ratchet: partial information



$F = F_{\text{stall}}$   
 $r$  infinite

Partial Information  
 $x = \{0, 1, 2\}$

If  $F = F_{\text{stall}}$ , the current vanishes but the system is out of equilibrium



Three-time correlations ( $d_3$ ) = A new footprint of irreversibility

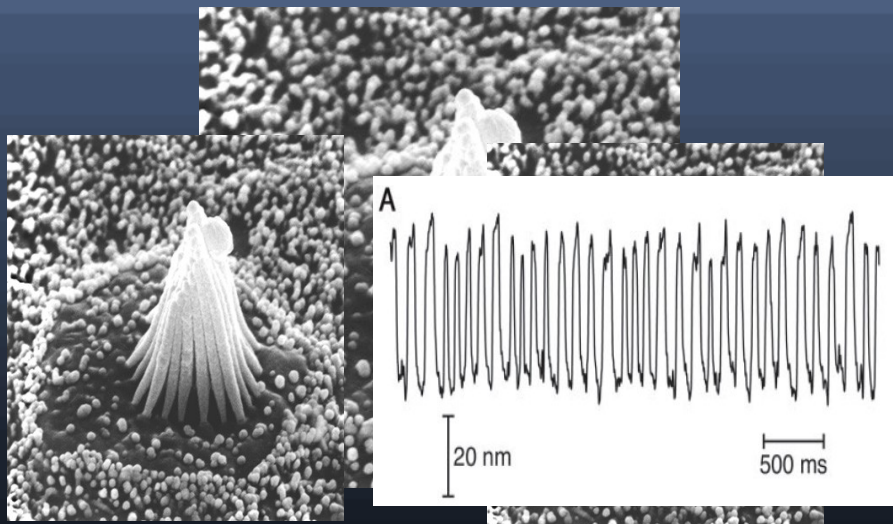
# Conclusion

- $d(p_F \parallel p_B)$  proved to be a good estimator of  $S_{\text{prod}}$  even ignoring physical details of the system
- Brute force estimators  $\ggg$  Ziv-Merhav estimator
- $F=0$  and  $r$  finite: The technique works qualitatively. Ansatz found.
- $F=0$  and  $r$  infinite: Markovianity is not enough
- $F=F_{\text{stall}}$  and  $r$  finite: Three-time correlations ( $d_3$ ) capture the irreversibility where the current is zero.

# Future work

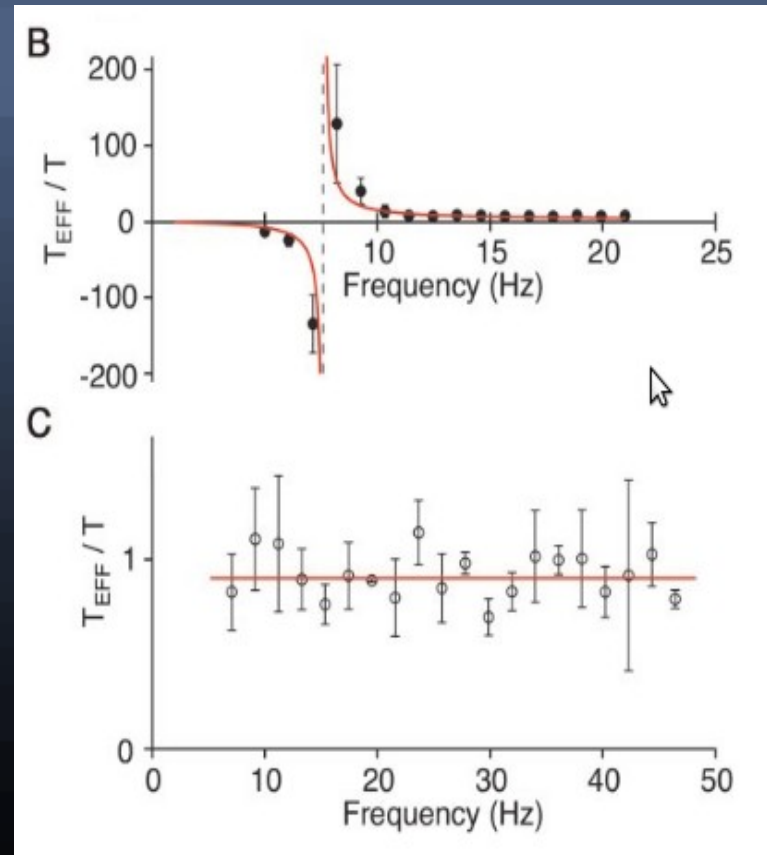
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Ear hair bundles



spontaneous  
oscillations

forced  
oscillations



Active cells

Passive cells



Thanks for your attention